

ORIGINAL ARTICLE

Fundamental mechanisms responsible for the temperature coefficient of resonant frequency in microwave dielectric ceramics

Shengke Zhang¹ | Hasan Sahin^{2,4} | Engin Torun² | Francois Peeters² | Dinesh Martien³ | Tyler DaPron³ | Neil Dilley³ | Nathan Newman¹

¹Arizona State University, Tempe, Arizona

²Department of Physics, University of Antwerp, Antwerp, Belgium

³Quantum Design, San Diego, California

⁴Department of Photonics, Izmir Institute of Technology, Izmir, Turkey

Correspondence

Nathan Newman, Arizona State University, Tempe, Arizona.

Email: Nathan.Newman@asu.edu

Abstract

The temperature coefficient of resonant frequency (τ_f) of a microwave resonator is determined by three materials parameters according to the following equation: $\tau_f = -(\frac{1}{2} \tau_\epsilon + \frac{1}{2} \tau_\mu + \alpha_L)$, where α_L , τ_ϵ , and τ_μ are defined as the linear temperature coefficients of the lattice constant, dielectric constant, and magnetic permeability, respectively. We have experimentally determined each of these parameters for Ba(Zn_{1/3}Ta_{2/3})O₃, 0.8 at.% Ni-doped Ba(Zn_{1/3}Ta_{2/3})O₃, and Ba(Ni_{1/3}Ta_{2/3})O₃ ceramics. These results, in combination with density functional theory calculations, have allowed us to develop a much improved understanding of the fundamental physical mechanisms responsible for the temperature coefficient of resonant frequency, τ_f .

KEYWORDS

density functional theory, dilatation/dilatometry, dopants/doping, electron spin resonance, microwave resonators

1 | INTRODUCTION

Low-loss temperature-compensated Ba(B'_{1/3}Ta_{2/3})O₃ (B' = Zn²⁺, Nb²⁺, Mg²⁺) perovskite dielectrics are used extensively in high-performance communication and Doppler systems. Of these, cellular base station filters have become the most common application.^{1,2} Such advanced microwave systems require that the relative temperature dependence of the dielectric's resonant frequency, $\tau_f = (1/f)(df/dT)$, be precisely set at zero or near-zero to achieve the required performance over a range of temperatures.³ A small non-zero τ_f is often chosen by system engineers to offset the small temperature dependence of the amplifiers. Despite the importance of τ_f for these practical applications, a strong first-principles understanding of what determines this important parameter has not been established. The goal of this work was to demonstrate that currently available density functional theory (DFT) calculations can be used to quantitatively model the physical processes involved in determining τ_f with reasonable accuracy.

Controlling τ_f to the desired value in low-loss materials has been achieved empirically by researchers and commercial manufactures by (1) doping, (2) alloying, and/or (3) combining multiphase composite ceramic materials.⁴⁻¹⁰ Surprisingly, at least to the authors of this manuscript, almost all of the commercial microwave dielectric products on the market today tune τ_f through the addition of varying amounts of magnetic additives, such as Ni, Co, and Mn.⁶⁻¹⁰ In this work, Ba(Zn_{1/3}Ta_{2/3})O₃ (BZT) was selected for the host material because it is the lowest-loss temperature-compensated microwave ceramic with a dielectric constant of ~30 or greater. 0.8% Ni-doping was used since this concentration of Ni doping in commercial materials tunes τ_f to near-zero.^{8,11} Finally, since the Ba(Zn_{1/3-x}Ni_xTa_{2/3})O₃ system is a completely miscible solid solution whose room-temperature dielectric constant and τ_f vary smoothly and monotonically over the entire alloy series, we believe that the inferences we derive from 0.8 at.% and Ba(Ni_{1/3}Ta_{2/3})O₃ (BNT) are general and can be used to understand other concentrations.^{6,9,10}

To develop a strong first-principles understanding, we first must determine how material properties are related to this important device parameter, τ_f . The resonant frequency, f , for a representative structure, such as a half-wavelength resonator of length l , is $\frac{c}{2l\sqrt{\epsilon\mu}}$ where c is the speed of light under vacuum, ϵ is the dielectric constant, and μ is the magnetic permeability.³ Then by taking partial derivatives with respect to temperature, we obtain the more general equations.

$$\begin{aligned} \frac{df}{dT} &= -\frac{c}{2l\sqrt{\epsilon\mu}} \times \frac{dl}{dT} - \frac{c}{4l\sqrt{\epsilon\mu}} \times \frac{d\epsilon}{dT} - \frac{c}{4l\sqrt{\epsilon\mu}} \times \frac{d\mu}{dT} \\ &= -f \times \frac{dl}{dT} - \frac{f}{2} \times \frac{d\epsilon}{dT} - \frac{f}{2} \times \frac{d\mu}{dT} \end{aligned} \quad (1)$$

$$\frac{df}{dT} = -\frac{dl}{dT} - \frac{1}{2} \times \frac{d\epsilon}{dT} - \frac{1}{2} \times \frac{d\mu}{dT} \quad (2)$$

Thus,

$$\tau_f = -\left(\frac{1}{2\tau_\epsilon} + \frac{1}{2\tau_\mu} + \alpha_L\right) \quad (3)$$

The three material parameters, $\alpha_L = (1/l)(dl/dT)$, $\tau_\epsilon = (1/\epsilon)(d\epsilon/dT)$, and $\tau_\mu = (1/\mu)(d\mu/dT)$ in Equation (3), are referred to the linear temperature coefficients of the lattice constant (ie, thermal expansion coefficient), dielectric constant, and magnetic permeability, respectively.^{1,2}

To better understand the mechanism responsible for these three parameters, it is useful to establish their temperature dependences from near 0 K, where thermal excitations of phonon, electronic and spin excitations are minimal, up to say their operating temperature of ~ 300 K. We are not aware of any published papers in the literature that have determined all the relevant parameters, τ_ϵ , τ_μ , α_L , and τ_f for a microwave ceramic material at cryogenic temperatures. In the next few paragraphs we will summarize the literature that investigated the material's properties of microwave perovskite's at low temperature.

There have been a handful of studies¹²⁻¹⁵ which have characterized only the parameter α_L for low-loss perovskite ceramics over this temperature range. In an investigation of the thermal expansion of $\text{Ba}(\text{Mg}_{1/3}\text{Ta}_{2/3})\text{O}_3$ ($\tau_f @ 300\text{K} = 5.4$ ppm/K), Gvasaliya et al.¹³ reported that the degree of $\text{Mg}^{2+}/\text{Ta}^{5+}$ does not significantly influence the material's thermal expansion. α_L of BaZrO_3 ($\tau_f @ 300\text{K} \sim 140$ ppm/K) was measured by neutron scattering at three temperatures between 2 and 300 K,^{14,15} making it difficult to observe systematic trends.

A number of other studies have measured τ_f using a TE_{018} or TE_{011} resonance mode of dielectric resonators¹⁶⁻²² and assumed that this is equal to τ_ϵ , neglecting the contributions from α_L and τ_μ . From the literature^{1,2,3} and the measurements reported here, we note that α_L can be on the

order of 5-10 ppm/K at 200 K and above for most microwave oxide ceramics,¹ whereas τ_μ can be of the same order at very low temperature in paramagnetic-laden dielectrics.²⁴ Since these References¹⁶⁻²² directly measured τ_f , we will only summarize their findings in that light. Subodh et al.¹⁹ found that as the temperature increases, the τ_f of V_2O_5 -doped Mg_2TiO_4 dielectrics is ~ -30 ppm/K at low temperatures and increases to ~ -50 ppm/K above 150 K. An investigation of V_2O_5 -doped Mg_2TiO_4 ceramics by Bhuyan et al.¹⁶ and single-crystal LaAlO_3 ²⁰ found that τ_f is positive at low temperature, increasing slowly as the temperature is raised to ~ 150 -200 K, then increasing more rapidly above 150-200 K. Similarly, Pamu et al. found that τ_f of CuO-doped $\text{Zr}_{0.8}\text{Sn}_{0.2}\text{TiO}_4$ is ~ 9 ppm/K from 10 to 250 K and then increases significantly up to room temperature.²¹ Likewise, Varma et al. reported that 0.5 at.% of a number of dopants including In, Ti, Mn, and Sb in $\text{Ba}(\text{Zn}_{1/3}\text{Ta}_{2/3})\text{O}_3$ ceramics have a τ_f of ~ -5 at 25 K that rises to +2 to 5 ppm/K at room temperature. Of these dopants in $\text{Ba}(\text{Zn}_{1/3}\text{Ta}_{2/3})\text{O}_3$, only the Ti dopant shows any substantial temperature dependence in the range from 200 to 300 K.²² In that study, higher doping levels were found to exhibit similar increases in τ_f with temperature, but with quantitatively different values over the entire range of temperatures. In contrast, Jacob et al.¹⁷ found that τ_f of CuO/ZnO doped $\text{Zr}_{0.8}\text{Sn}_{0.2}\text{TiO}_4$ is essentially constant at ~ 5 ppm/K from 13 to 300 K. Interestingly, Sudheendran et al.¹⁸ presented that at cryogenic temperatures, τ_f of 0.9 at.% Ti^{4+} -doped $\text{Bi}_2(\text{Zn}_{2/3}\text{Nb}_{4/3})\text{O}_7$ could be negative from ~ 10 to 75 K; then it becomes positive and almost constant at the higher temperatures. In an early investigation, Tamura also observed that τ_f was nonmonotonic with temperature.²⁵ From this, it is clear that the temperature dependence of τ_f varies widely and an improved understanding of the mechanisms that determine this important property is needed.

In our study, we have experimentally measured α_L , τ_f , and τ_μ of BZT, 0.8 at.% Ni-doped BZT and $\text{Ba}(\text{Ni}_{1/3}\text{Ta}_{2/3})\text{O}_3$ (BNT), allowing us to determine τ_ϵ over the temperature range 2-300 K. We also performed first-principles calculations for the electronic and phonon structures of the compounds. This has allowed us to gain a much improved fundamental quantitative understanding of the mechanisms responsible for the observed temperature dependence of these parameters.

2 | EXPERIMENTAL PROCEDURE

$\text{Ba}(\text{Zn}_{1/3}\text{Ta}_{2/3})\text{O}_3$ (BZT), $\text{Ba}(\text{Ni}_{1/3}\text{Ta}_{2/3})\text{O}_3$ (BNT), and 0.8 at.% Ni-doped BZT dielectrics studied in this work were synthesized by conventional ceramic powder processing methods in our laboratory.²³ The resonant frequencies were determined using microwave dielectric resonator (DR)

measurements. The 0.85 cm × 0.60 cm × 0.20 cm rectangular DRs were inserted into a 2.03 cm × 0.61 cm × 1.52 cm gold-plated copper cavity which was affixed to the end of a cryogenic dipping probe. A HP8510C microwave vector network analyzer (Hewlett-Packard, Palo Alto, CA) was used to measure S_{21} vector transmission values. Micrometers are used to adjust the distance between the microwave coupling probes and sample that is weakly coupled to the resonator, allowing the unloaded Q of the DR to be determined to better than a few percent and the corresponding resonant frequency to be determined to typically better than 5 ppm. The microwave cavity was lowered into a cryostat (PPMS, Model 6000; Quantum Design, San Diego, CA) and the resonant frequency of the $TE_{11\delta}$ mode of rectangular DRs was measured in the range 9.5–10.0 GHz.^{23,24} Static dielectric constants, ϵ_0 , were measured by a conventional $TE_{01\delta}$ measurements technique on ~2.7 cm diameter, ~0.5 cm tall cylindrical samples in an 8.0 cm diameter × 7.0 high cylindrical gold-plated copper cavity at room temperature.^{23,24}

Magnetic permeability determinations of $\mu=1 + M/H$ are performed using a vibrating sample magnetometer option for the Quantum Design PPMS system (Model 6000; Quantum Design, San Diego, CA).

For this study, we used a commercial direct-reading capacitance dilatometer from Quantum Design (Model PPMS with the dilatometer option). This design directly infers temperature-induced changes in the sample length from the reactance difference between two capacitors in the dilatometer cell. In this regard, an absolute capacitance measurement is not required and the resulting symmetrical configuration reduces the effects of absorbed gasses and temperature gradients. The cell is constructed from fused silica, resulting in a low thermal expansion background.

These methods were chosen to determine τ_f , τ_e , τ_μ , and α_L with the required precision of better than a ppm.^{26,27} Ideally, a direct method to determine the dielectric constant as a function of temperature would be preferred, but unfortunately the established methods for this, such as the coaxial probe, transmission line, or free space transmission or reflection techniques,²⁸ are only accurate to a few percent under typical measurement conditions and to a little better than a part per thousand under the most optimal laboratory conditions.^{26,27} Nonresonant indirect (device) techniques, such as measuring the capacitance in the parallel plate geometry, do not have any better precision than a part per thousand too and would also need to be corrected for lattice expansion.²⁶⁻²⁸ Thus, neither the direct methods nor nonresonant indirect methods offer the precision needed a few ppm/K needed for this study. So, we used an electromagnetic resonator technique with the required precision to infer τ_e by correcting the measured temperature-dependent frequency shifts (τ_f) with suitably accurate measurements

of the lattice expansion (α_L) and the temperature dependence of magnetic permeability (τ_μ) (ie, Equation 3).

The DFT calculations were performed on BZT and BNT crystal structures using plane wave projector-augmented wave method²⁹ which is implemented in the Vienna Ab-initio Simulation Package³⁰⁻³² The generalized gradient approximation within the Perdew-Burke-Ernzerhof³³ is employed to describe the exchange-correlation potential. The kinetic energy cutoff for plane-wave expansion is set to 500 eV and the hexagonal Brillouin Zone was sampled using 14 × 14 × 12 Γ centered k-point grid for the primitive cell and for the examination of magnetic ground state it is scaled according to size of the supercell. The convergence criterion for energy is set to 10⁻⁵ eV between two consecutive steps in the self-consistent field calculations. Full structural relaxation was performed until the Hellmann-Feynman forces are less than 10⁻⁴ eV/Å. Phonon calculations were performed using small displacement methodology implemented in the PHON code.³⁴ The imaginary part of the electronic contribution of the dielectric function of BZT and BNT is calculated by summing over empty electronic states using Equation (4).

$$\epsilon_{\alpha\beta}(w) = \frac{4\pi^2 e^2}{\Omega} \lim_{q \rightarrow 0} \frac{1}{q^2} \sum_{c,v,k} 2w_k \delta(\epsilon_{ck} - \epsilon_{vk} - w) \times \langle u_{ck+e_{\alpha q}} | u_{vk} u_{vk} | u_{ck+e_{\beta q}} \rangle \quad (4)$$

where q , Ω , ϵ , w , u_{ck} , c , and v correspond to Bloch vector, volume of the unit cell, dielectric function, frequency, cell periodic part of the orbitals at the k -point (k), conduction, and valence band states, respectively.

For this work, we only calculated the electronic and phonon band structures of the compounds Ba(Zn_{1/3}Ta_{2/3})O₃ and Ba(Ni_{1/3}Ta_{2/3})O₃ and did not carry out simulations for the random alloy Ba([Zn_{1-x}Ni_x]_{1/3}Ta_{2/3})O₃ due to the very expensive computational cost required to accurately simulate a very large number of atoms required for random alloys with exchange interactions.

3 | RESULTS AND DISCUSSION

In Figure 1, we present our experimental measurements of τ_f , the factor of interest to microwave system designers. Then, in the remainder of the section, we give experimental and simulation results for α_L , τ_e , and τ_μ , parameters that describe the temperature dependence of the physical properties that determine this important factor.

3.1 | Determining α_L

Next, we provide experimental determinations of α_L and show that it can be quantitatively modeled and understood

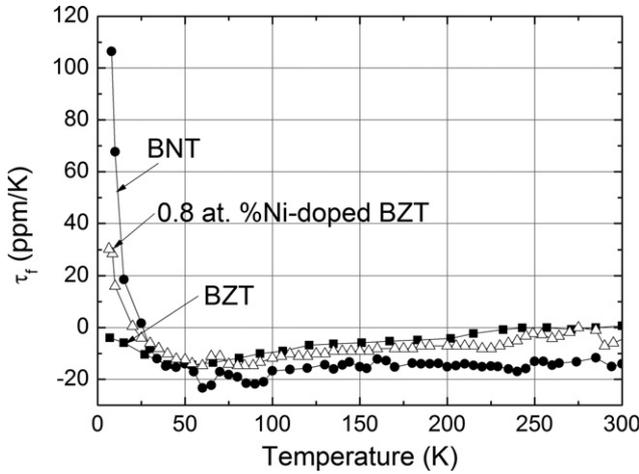


FIGURE 1 Temperature dependence of temperature coefficient of resonant frequency (τ_f) in BZT, BNT, and 0.8 at.% Ni-doped BZT.

using our DFT calculations. In Figure 2A, we see that the thermal BZT and BNT expansion data up to ~ 200 K can be accurately fit with the Debye equation for the specific heat and Grüneisen analysis³⁶ using a characteristic Debye temperature of 392 and 408 K, respectively. When the temperature is above 200 K, α_L starts to deviate from the Grüneisen model as a result of the extra weight of softened modes near the Brillouin zone boundaries due to an antiferromagnetic transition (Figure 3). It is also noticed that the thermal expansion reduces as the Ni concentration increases, which is often correlated with the smaller atomic weight and shorter effective ionic radii of the Ni^{2+} ion than Zn^{2+} .³⁸

The results of the DFT calculations of the phonon dynamics are shown in Figure 3. From the slope of the

lowest energy acoustic phonon mode of the simulation, we estimate the Debye temperature of BZT and BNT to be ~ 438 K and ~ 480 K, respectively, $\sim 15\%$ higher than found when fitting experimental data. This is a reasonable agreement, given that our estimate of the Debye temperature from the theory comes only from only the lowest energy acoustic mode. Both theory and experiment find that BNT has a larger Debye temperature than BZT. Given that Ni has a lighter atomic mass than Zn by $\sim 10\%$, we would expect BNT to have a higher sound velocity and thus Debye temperature. Even though simple models indicate that these properties scale with the effective spring constant and inversely with the square root of the atomic mass, it is difficult to make quantitative estimates given that the Debye temperature in these materials is established by many vibrational modes each involving a number of the compound's elemental constituents.

Using Grüneisen analysis, we can infer α_L from the DFT calculations of the heat capacity (C_v) using Equation (5). Results for α_L as a function of temperature are shown in Figure 2B.

$$\alpha = \frac{\gamma C_v}{3B} \quad (5)$$

where γ is the average Grüneisen parameters and B is the bulk modulus.³⁰ We calculate the Grüneisen parameter (γ) of BZT and BNT using, $\gamma = -\frac{d \ln(\omega)}{d \ln(V)}$ where V and ω correspond to volume of the unit cell and the average Debye frequency of the acoustic modes. The Grüneisen parameter of BZT and BNT are calculated to be 2.39 and 1.66, respectively. This analysis indicates that small thermal expansion coefficients are expected from compressible

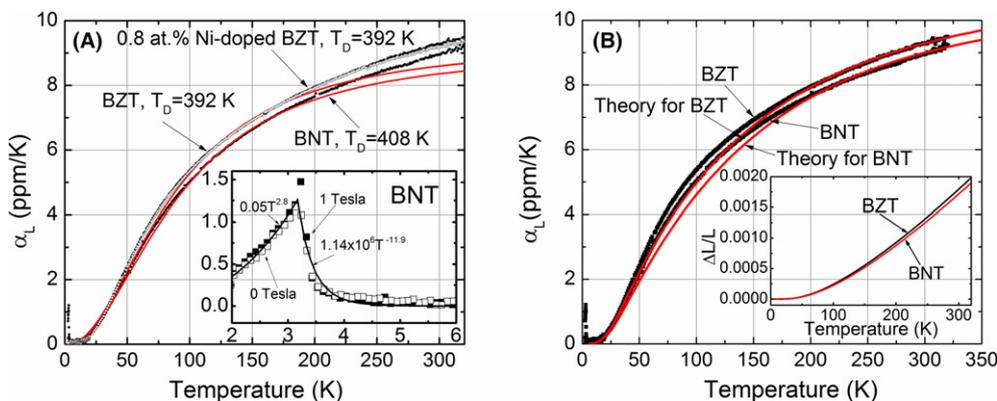


FIGURE 2 A, Thermal expansion (α_L) measurement values for BZT and BNT are represented as black data points and the Grüneisen model fits are represented as red solid lines. The inset shows BNT's Neel temperature at 3.3 K, characteristic of a phase transition from the antiferromagnetic state at low temperatures to the paramagnetic state.³⁵ A peak in the BNT magnetic permeability at 3.3 K (inset of Figure 4B) confirms this. Neither BZT or 0.8 at.% Ni-doped BZT show evidence for a peak in their thermal expansion characteristics, indicating that they are not antiferromagnetic over the measured temperature range. B, α_L (black data points) and density functional theory results (red solid lines) of BZT and BNT. The inset shows the temperature dependence of lattice expansion of BZT and BNT, as derived from the integration of the experimental α_L . [Color figure can be viewed at wileyonlinelibrary.com]

lattices whose high-energy phonons do not significantly soften upon expansion (ie, small Grüneisen parameters). Thus, materials with high Debye temperatures within a given class of material will tend to have smaller thermal expansion coefficients, as we found when comparing the results for BZT to BNT.

The values of $B_{\text{BZT}}=209$ GPa and $B_{\text{BNT}}=152$ GPa were inferred by scaling the theory relation to the experiment values in Figure 2B. We note that the theoretical results closely follow the measurement values, particularly below 50 K and above 175 K. These observations again confirm that the theoretical calculations can adequately model the phonon dynamics in these materials.

We also used the theory alone to model the bulk modulus of BZT and BNT. We changed the volume of the unit cell from -2% to 2% with step of 1% and obtained the

energy (E) as a function of volume (V). Then, a second order polynomial function was used to fit the E vs V data, and the bulk moduli of BZT and BNT were calculated by taking the second order derivative of E with respect to the V , as indicated by Equation (6).

$$B = V \frac{d^2 E}{d^2 V} \quad (6)$$

where B and E correspond to the bulk modulus and energy, respectively. Our calculations show that the bulk moduli of BZT and BNT are 138.10 and 138.75 GPa, respectively. These values are significantly different than inferred from the Grüneisen analysis of the experimental data reported earlier in this section. The variance may arise as a result of using an average Grüneisen parameter rather than

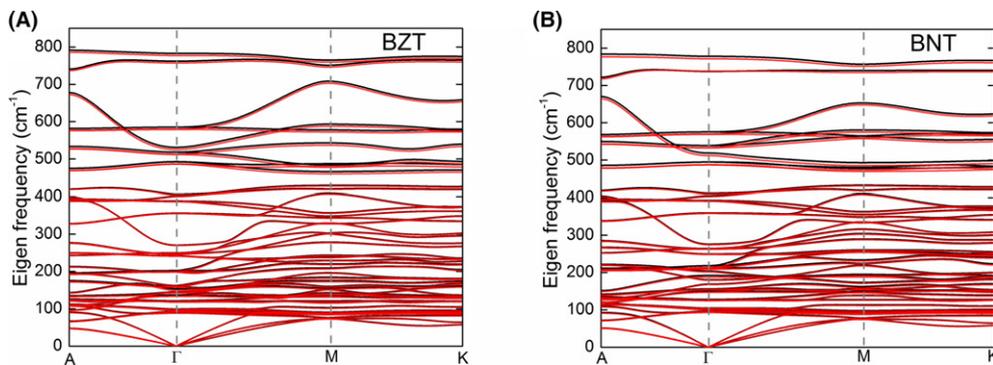


FIGURE 3 Phonon Eigen frequency ($\omega=E/h$) vs wave vector (k) curves of BZT (A) and BNT (B) with relaxed lattice at equilibrium and at 0 K (black lines) and at a 0.2% lattice enlargement (red lines), the experimentally measured increase at ~ 300 K (inset Figure 2B). The 0.2% expanded lattice is used to simulate the properties at 300 K. This is based on the concept that a solid's properties can be adequately modeled by considering the lattice softening alone upon lattice expansion, as has been generally accepted by the field³⁷ [Color figure can be viewed at wileyonlinelibrary.com]

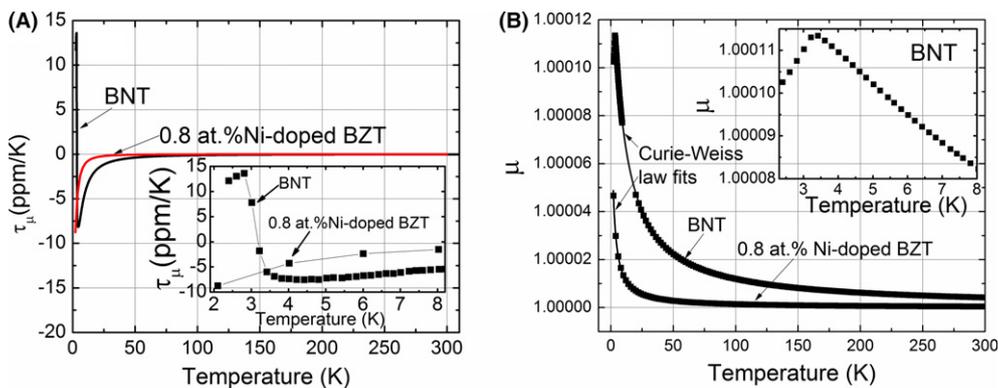


FIGURE 4 (A) Temperature coefficient of magnetic permeability (τ_{μ}) of BNT and 0.8 at.% Ni-doped BZT characterized measured using a vibrating sample magnetometer at 40 Hz. The inset shows the τ_{μ} from 2 to 8 K; (B) temperature dependence of permeability (μ) of BNT and 0.8 at.% Ni-doped BZT, and black solid lines are the Curie-Weiss law fits. The inset shows the antiferromagnetic Neel temperature of BNT is 3.3 K [Color figure can be viewed at wileyonlinelibrary.com]

performing a k -dependent analysis for every phonon mode. More detailed analysis will be reported in the future.

3.2 | Determining τ_μ

Next we show the temperature dependence of permeability (μ) and its temperature-derivative, τ_μ , in Figure 4. Note that the μ of BNT follows the expected Curie-Weiss law in Equations (7) and (8).³⁶

$$\mu = \frac{C}{(T + \Theta)} \quad (7)$$

$$\tau_\mu = -\frac{C}{(T + \Theta)^2} \quad (8)$$

where Θ is Weiss temperature and C is the Curie constant. The Θ of BNT is calculated to be -8.37 K. When the temperature is below the Neel temperature, 3.3 K,

the BNT transitioned to antiferromagnetic phase, which shows a discontinuity in both μ and τ_μ . The magnitude of the Weiss temperature is a little under three times that of the Neel temperature, indicating that the antiferromagnetically ordered state is established by interactions that extend beyond just the effective fields of the nearest neighbors, to the second nearest neighbor and potentially beyond.³⁶

The magnetic susceptibility data as a function of temperature of the 0.8 at.% of Ni-doped BZT can be fit to the Curie-Weiss law (Equation 6) with a Weiss temperature, Θ , of -1.12 K, as shown in Figure 4B. Since the Weiss temperature is non-zero and we did not observe evidence for a phase transition during our measurements, this suggests that this level of doping will be antiferromagnetic at a temperature below the lowest temperature measured. In contrast, the undoped BZT magnetic susceptibility data

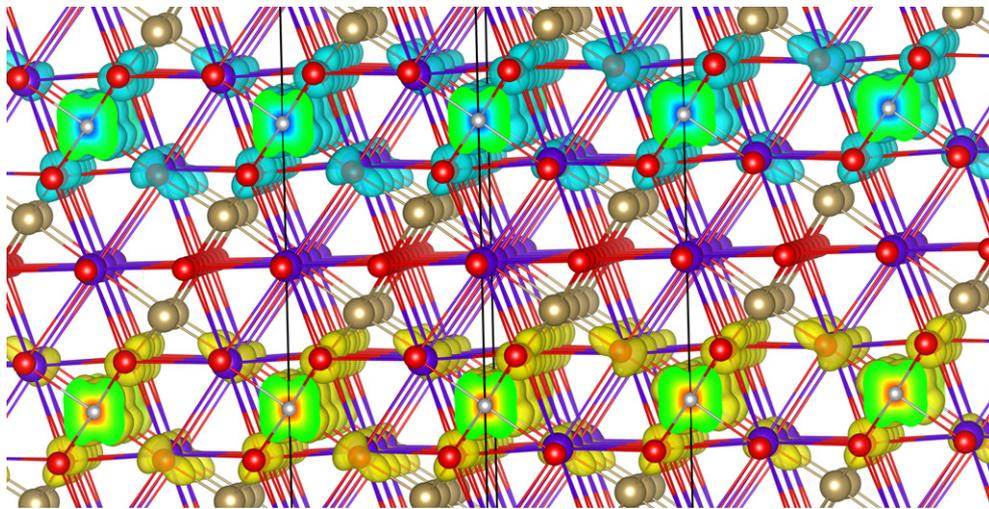


FIGURE 5 The magnetization charge density plot of antiferromagnetic structure of BNT crystal. The Ba, Ta, Ni, and O atoms are shown in purple, gold, gray, and red, respectively. The majority and the minority magnetization densities are shown in yellow and blue, respectively. The isosurface value is set to $0.003 \text{ e}/\text{\AA}^3$ [Color figure can be viewed at wileyonlinelibrary.com]

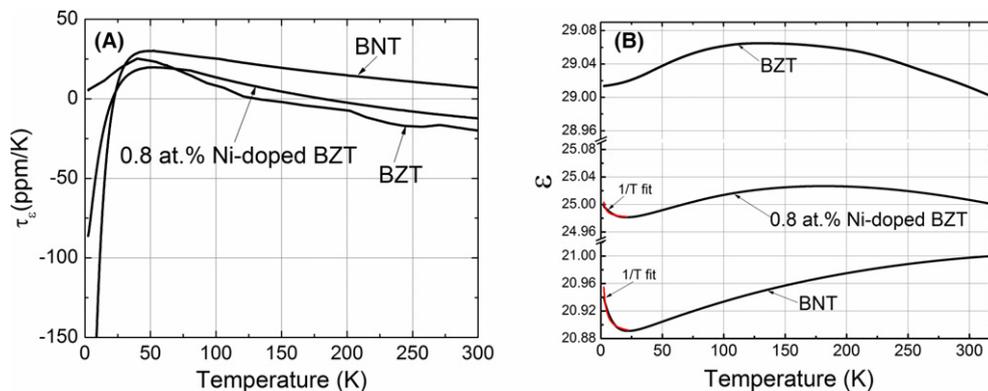


FIGURE 6 (A) The temperature coefficients of the dielectric constant (τ_ϵ) for BZT, BNT, and 0.8 at.% Ni-BZT are inferred from τ_f , α_L , and τ_μ ; (B) the dielectric constant (ϵ) is obtained by integrating the τ_ϵ and then multiplying the normalized dielectric constants with the ϵ_0 of 29, 21, and 25 for BZT, BNT, and 0.8 at.% Ni-BZT, respectively. The red solid lines at low temperatures are fits to the Curie law which shows reasonable agreement. [Color figure can be viewed at wileyonlinelibrary.com]

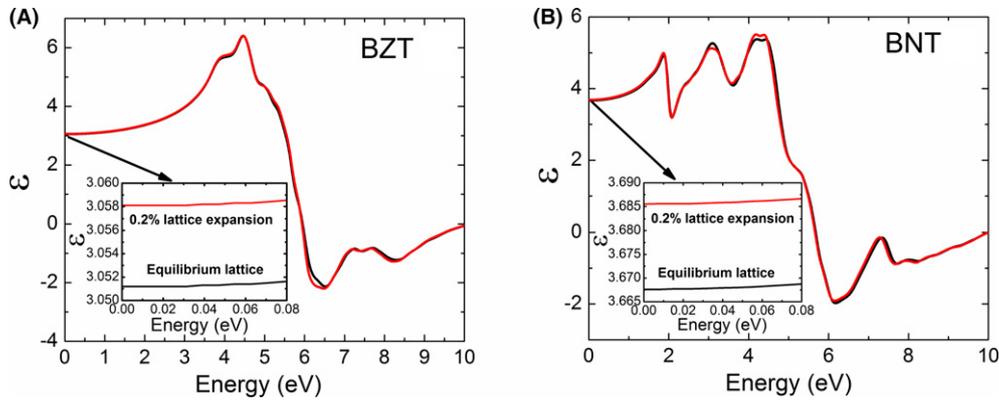


FIGURE 7 The electronic contribution to the dielectric constant of BZT (A) and BNT (B) with relaxed lattice and 0.2% of lattice expansion from density functional theory calculations [Color figure can be viewed at wileyonlinelibrary.com]

were found to be diamagnetic from 2 to 300 K and thus it does not fit the Curie law in this temperature range.

DFT calculations show that in BNT the ground-state antiferromagnetic configuration at 0 K exists in an intra-layer-ferromagnetic and interlayer-antiferromagnetic structure as shown in Figure 5. Figure 5 shows that the magnetization density is not only localized on the Ni atoms. The nearest-neighbor O atoms of Ni atoms also carry magnetization and that is why the structure is intra-layer-ferromagnetic. This configuration is 3.2 meV per cell lower in energy than the ferromagnetic configuration.

3.3 | Determining τ_ε

To obtain the last unknown physical parameter, τ_ε , the temperature dependence of the dielectric constant (τ_ε) can be inferred from $\tau_\varepsilon = -(2\tau_f + 2\alpha_L + \tau_\mu)$, as shown in Figure 6A. The red solid lines in Figure 6B show that from 2 to 20 K, the magnitude of ε of 0.8 at.% Ni-BZT and BNT roughly scales with the Ni content and can be reasonably fit with a Curie dependence with equation of $\varepsilon_\infty + C/T$, where ε_∞ is the baseline dielectric constant before it changed its slope, a relation that is found for paraelectric materials or materials with paraelectric defects.^{37,39-41} Since the DFT calculations of the phonon characteristics do not show any indication of ferroelectricity in the materials studied, this observation must arise from the presence of paraelectric defects associated with Ni. This can result from reorientation over an energy barrier of an asymmetric (1) Ni atom in the lattice, (2) spectator defect around the Ni, or (3) d-shell electronic state. This is a surprising result and warrants further study.

We will now show why the temperature dependence of dielectric constant in BNT is different from the pure BZT compound, and how Ni influences the phonon spectrum and its associated dielectric constant shift with temperatures. As can be determined by the DFT calculations of the electronic contribution to ε_r in Figure 7, BZT's electronic

contribution to the microwave dielectric constant (ie, frequencies of ~ 0.5 -100 GHz) increases by 0.007 when it undergoes a 0.2% lattice expansion. For BNT, ε_r is increased by ~ 0.02 from 3.667 to 3.685. Thus, the changes in the electronic polarizability accounts for $\sim 70\%$ (0.01) and $\sim 33\%$ (0.06) of the shift associated with ε_r of BZT and BNT over this range.

In future studies, we hope to extend current DFT methods of calculating the phonon contribution to the dielectric constant⁴² to model its temperature dependence.

4 | CONCLUSION

In this study, we used experimental methods to determine τ_f , τ_ε , τ_μ , and α_L over the temperature range of 2-300 K for BZT, 0.8 at.% Ni-doped BZT, and BNT. These results, in combination with DFT calculations, have allowed us to develop a much improved understanding of the fundamental mechanisms responsible for determining τ_f for these high-performance microwave dielectrics. From our work, we concluded:

- 1 A Grüneisen analysis using the Debye specific heat model can quantitatively fit the thermal expansion data over the entire temperature range investigated.
- 2 From 2 to 300 K, 70% of the observed temperature coefficient of resonant frequency (τ_f) arises from the electronic contribution to the dielectric constant in BZT, in contrast to BNT where virtually the same fraction arises from the phonon contribution to the dielectric constant.
- 3 At temperatures below 25 K, a Curie dependence from paraelectric defects dominates the temperature coefficient of resonant frequency (τ_f).
- 4 The magnetic contribution to the temperature coefficient of resonant frequency (τ_f) in Ni-containing material arises from spin excitations, but is over an order of magnitude smaller than the other contributions over all temperatures measured.

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